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# Success Probabilities for Linear Optics Gates



Institute for  

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Mathematical Sciences

Koenraad M.R. Audenaert



Quantum Information at  
Imperial College  
London



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# Impressum

Extended Club Remix of work by Scheel, Lütkenhaus, Audenaert, Eisert:

- S. Scheel, *quant-ph / 0406127*.
- S. Scheel and N. Lütkenhaus, NJP **6**, 51 (2004).
- S. Scheel and K. Audenaert, NJP **7**, 149 (2005).
- J. Eisert, PRL **95**, 040502 (2005).

Unified presentation, shorter proofs, more general results, and a small ToDo-list...



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# A basic component in linear-optics quantum computing

- One of the promising routes to implementing small-scale quantum networks is by using linear-optical networks.
- Certain two-mode gates (e.g. controlled- $\sigma_z$ ) can be generated by acting separately on both modes within a Mach–Zehnder interferometric setup.
- In that way the complexity of a two-mode gate is reduced to a single-mode gate: the nonlinear sign-shift gate

$$c_0|0\rangle + c_1|1\rangle + c_2|2\rangle \mapsto c_0|0\rangle + c_1|1\rangle - c_2|2\rangle$$

- A number of theoretical works (Ralph, Scheel, Lapaire, KLM, Lufthansa) give a variety of networks capable of implementing this gate.



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# Subject of the talk

- The crucial question for the ability to concatenate linear-optical gates are their probabilities of success.
- In these works we studied the question how the probability of success scales as the signal state becomes higher-dimensional, thereby giving a hint on the possible scaling law for multi-mode quantum gates.
- We considered the  $N + 1$ -dimensional sign-shift gate

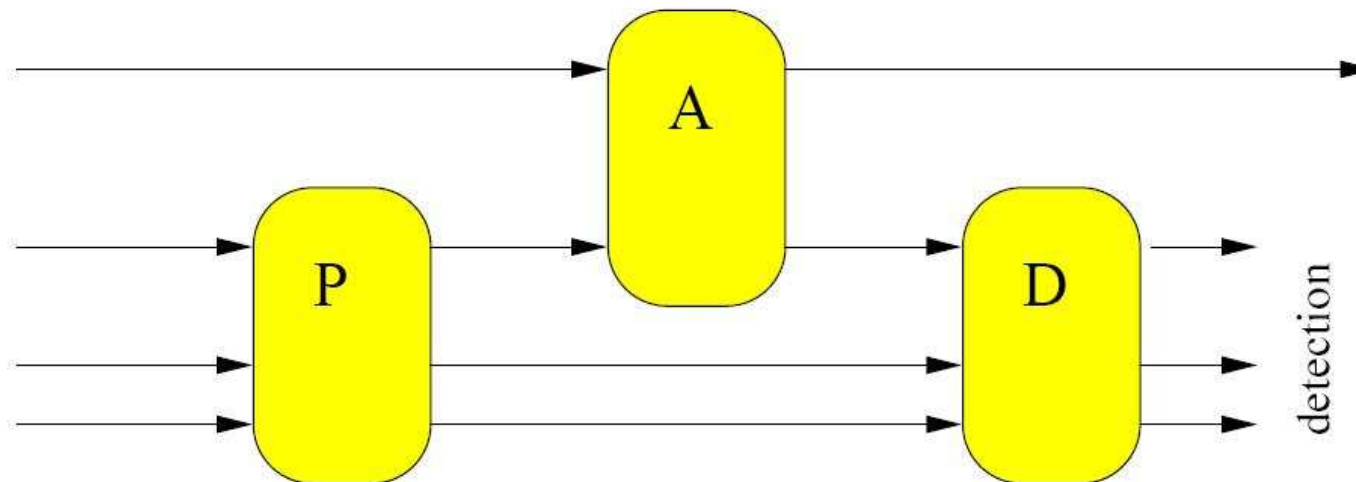
$$c_0|0\rangle + \dots + c_{N-1}|N-1\rangle + c_N|N\rangle \mapsto c_0|0\rangle + \dots + c_{N-1}|N-1\rangle - c_N|N\rangle.$$

- We limit ourselves to single-shot implementations without conditional feed-forward.
- We consider an abstract three-stage type of network that includes all possible networks.



# Basic Circuit

- A: active beam splitter
- P: ancilla preparation
- D: ancilla projection





# Basic Circuit

- The generality of this abstraction follows from the observation that each  $U(N)$  network can be decomposed into a triangle-shaped network of at most  $N(N - 1)/2$  beam splitters (and some additional phase shifters) [Reck et al, 94].
- Then, the signal mode can be chosen to impinge only on a single beam splitter ‘A’, with the rest of the network divided into the ancilla preparation and detection stages ‘P’ and ‘D’ [Scheel and Lütkenhaus, 04].
- The ‘D’ stage is a two-outcome POVM:
  - outcome ‘1’ = “accept”;
  - outcome ‘0’ = “bad luck; try again”
- We can show that having just one “accept” outcome is optimal.



# Basic Circuit

- A: active beam splitter,  $\hat{U}$ , transmittivity  $T$ , reflectivity  $R$ ;  $|T|^2 + |R|^2 = 1$
- P: ancilla preparation,  $|\phi\rangle = \sum_{l=0}^n \gamma_l |l\rangle|n-l\rangle$
- D: ancilla projection on  $|\psi\rangle = \sum_{l=0}^n \alpha_l^* |l\rangle|n-l\rangle$
- The vectors  $\alpha$  and  $\gamma$  are, of course, normalised.
- The states  $|n-l\rangle$  can be multimode, and play no further role.
- The matrix elements of  $\hat{U}$  in the  $n$ -photon sector are

$$\langle m_1, m_2 | \hat{U} | n_1, n_2 \rangle = \frac{1}{\sqrt{m_1! m_2! n_1! n_2!}} \text{Per} \begin{pmatrix} TE_{m_1, n_1} & RE_{m_1, n_2} \\ -R^* E_{m_2, n_1} & T^* E_{m_2, n_2} \end{pmatrix},$$

with  $m_1 + m_2 = n_1 + n_2 = n$ , and  $E_{j,k}$  denotes a  $j \times k$  matrix all of whose entries are 1.



# Basic Circuit Behaviour

- Let the input state be the basis vector  $|k\rangle$ .
- By orthogonality of the states  $|n-l\rangle$ , the post-measurement state  $|\psi_k\rangle$  is

$$|\psi_k\rangle = \sum_{j=0}^N \sum_{l=0}^n \alpha_l \gamma_l \langle j, l | \hat{U} | k, l \rangle |j\rangle.$$

- Since a lossless beam splitter is photon-number preserving, there are only contributions from the term  $j = k$ .
- Denote  $\alpha_{kl} := \langle k, l | \hat{U} | k, l \rangle$ , then

$$|\psi_k\rangle = \left( \sum_{l=0}^n \alpha_l \gamma_l \alpha_{kl} \right) |k\rangle.$$

- As the states  $|n-l\rangle$  have left the picture, we can henceforth let  $n$  be as large as imaginable, say  $\infty$ .



# A Formula for the Success Probability

- Let  $A$  be the matrix with entries  $a_{jk}$ ,  $j = 0 \dots N$  and  $k = 0 \dots \infty$ .
- **Theorem** [Scheel and Audenaert]: The maximal success probability for obtaining from  $A$  the non-linear gate  $c_k \mapsto y_k c_k$ , for  $k = 0 \dots N$ , is given by

$$P_{\max} = \max_x \left\{ \frac{1}{|x|_1^2} : Ax = y \right\}$$

- $|\cdot|_1$  is the  $\ell_1$  vector norm.
- The minimisation is over all infinite-dimensional complex vectors  $x = (x_j)_{j=0}^{\infty}$ .



# Proof

- The transformation realised using ancilla  $\gamma_l$  and measurement  $\alpha_l^*$  is

$$c_k \mapsto y'_k c_k \quad (k = 0 \dots N), \text{ with } y'_k = \sum_{l=0}^{\infty} a_{kl} (\alpha_l \gamma_l).$$

- More concisely, I write this as  $y' = A(\alpha\gamma)$ , where  $A = (a_{kl})_{k=0, l=0}^{N, \infty}$ .
- The target transformation is realised, with probability  $P$ , if and only if

$$y' = \sqrt{P} y.$$

- To find a realisation with maximal  $P$ , first fix a vector  $x$  s.t.  $y = Ax$ .
- Find normalised  $\alpha$  and  $\gamma$  s.t.  $\alpha\gamma = \sqrt{P}x$  and  $P$  is as large as possible.



# Proof

- Find normalised  $\alpha$  and  $\gamma$  s.t.  $\alpha\gamma = \sqrt{P}x$  and  $P$  is as large as possible.
- Fix  $\gamma$ . Then  $\alpha = \sqrt{P}x/\gamma$ . Since  $|\alpha|_2$  must be 1, this fixes  $P = 1/|x/\gamma|_2^2$ .
- Now maximise this  $P$  by varying over all normalised  $\gamma$ .
- We find  $\min_{\gamma}\{|x/\gamma|_2 : |\gamma|_2 = 1\} = |x|_1$ .

Proof: using Cauchy-Schwarz

$$|x/\gamma|_2 |\gamma|_2 \geq \sum_l |(x_l/\gamma_l)\gamma_l| = |x|_1,$$

and noting that equality can be obtained (with  $|\gamma_l|^2 = c|x_l|$ ).

- Hence the maximum  $P$  for given  $x$  (s.t.  $Ax = y$ ) is  $P = 1/|x|_1^2$ .
- Finally, we maximise this by varying over all  $x$  s.t.  $Ax = y$ . □



# A Formula for the Success Probability

- **Corollary** [Scheel and Audenaert]: The maximal success probability for obtaining from  $A$  the non-linear gate  $c_k \mapsto y_k c_k$ , for  $k = 0 \dots N$ , is given by

$$P_{\max} = \max_x \left\{ \frac{1}{|x|_1^2} : Ax = y \right\}$$

where the  $x$  can be restricted to those that have at most  $N + 1$  non-zero entries.

- Geometrical proof



# Numerical method

- Corollary reduces the optimisation over  $\mathbb{C}^\infty$  to an optimisation over  $\mathbb{N}^{N+1}$ .
- We can replace  $x$  by a list of  $N + 1$  **positions**,  $0 \leq n_1 < n_2 < \dots < n_l$ , and corresponding **values**  $x_l$ .
- The optimisation over  $x_l$  can be solved directly by matrix inversion: only one  $x$  is feasible, namely  $x^0 := \tilde{A}^{-1}y$ . Here,  $\tilde{A}$  is an  $(N + 1) \times (N + 1)$  matrix containing the  $N + 1$  columns of  $A$  designated by the positions  $n_l$ .
- While this observation dramatically simplifies numerical approaches, the remaining optimisation over the  $n_l$  still makes for tedious business if we are after analytical results.
- For that reason we will consider an alternate expression for the maximal success probability, first obtained by Jens Eisert.



# Upper bound on Success Probability

- Let  $A$  again be the matrix with entries  $a_{jk}$ ,  $j = 0 \dots N$  and  $k = 0 \dots \infty$ ,
- **Theorem** [Eisert]: The maximal success probability for obtaining from  $A$  the non-linear phase-shift gate  $c_k \mapsto y_k c_k$ , with  $|y_k| = 1$ , for  $k = 0 \dots N$ , is upper bounded by

$$P_{\max} \leq \min_s \left( \frac{|A^\dagger(y\mathbf{s})|_\infty}{|\sum_{j=0}^N s_j|} \right)^2.$$

- $|\cdot|_\infty$  is the  $\ell_\infty$  vector norm,
- The minimisation is over all  $(N + 1)$ -dimensional (finite!) complex vectors  $\mathbf{s} = (s_j)_{j=0}^N$ .



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# Proof

- Hölder's inequality for vectors:  $|\langle a|b\rangle| \leq |a|_1 |b|_\infty$ .



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# Proof

- Hölder's inequality for vectors:  $|\langle a|b\rangle| \leq |a|_1 |b|_\infty$ .
- In our case: put  $a = x$  and  $b = A^\dagger(ys)$

$$|A^\dagger(ys)|_\infty |x|_1 \geq |x^\dagger A^\dagger(ys)|$$



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# Proof

- Hölder's inequality for vectors:  $|\langle a|b\rangle| \leq |a|_1 |b|_\infty$ .
- In our case:  $Ax = y$

$$|A^\dagger(y_s)|_\infty |x|_1 \geq |x^\dagger A^\dagger(y_s)| = |y^\dagger(y_s)|$$



# Proof

- Hölder's inequality for vectors:  $|\langle a|b\rangle| \leq |a|_1 |b|_\infty$ .
- In our case:

$$|A^\dagger(y_s)|_\infty |x|_1 \geq |x^\dagger A^\dagger(y_s)| = |y^\dagger(y_s)| = \left| \sum_j s_j |y_j|^2 \right|$$



# Proof

- Hölder's inequality for vectors:  $|\langle a|b\rangle| \leq |a|_1 |b|_\infty$ .
- In our case:  $|y_j| = 1$

$$|A^\dagger(y s)|_\infty |x|_1 \geq |x^\dagger A^\dagger(y s)| = |y^\dagger(y s)| = \left| \sum_j s_j |y_j|^2 \right| = \left| \sum_j s_j \right|$$



# Proof

- Hölder's inequality for vectors:  $|\langle a|b\rangle| \leq |a|_1 |b|_\infty$ .
- In our case:

$$|A^\dagger(y_s)|_\infty |x|_1 \geq |x^\dagger A^\dagger(y_s)| = |y^\dagger(y_s)| = \left| \sum_j s_j |y_j|^2 \right| = \left| \sum_j s_j \right|$$

- Thus, for all  $s$  and for all  $x$  satisfying  $Ax = y$ :

$$\frac{|A^\dagger(y_s)|_\infty}{\left| \sum_{j=0}^N s_j \right|} \geq \frac{1}{|x|_1},$$



# Proof

- Hölder's inequality for vectors:  $|\langle a|b\rangle| \leq |a|_1 |b|_\infty$ .
- In our case:

$$|A^\dagger(y)s|_\infty |x|_1 \geq |x^\dagger A^\dagger(y)s| = |y^\dagger(y)s| = \left| \sum_j s_j |y_j|^2 \right| = \left| \sum_j s_j \right|$$

- Thus, for all  $s$  and for all  $x$  satisfying  $Ax = y$ :

$$\frac{|A^\dagger(y)s|_\infty}{\left| \sum_{j=0}^N s_j \right|} \geq \frac{1}{|x|_1},$$

- Minimise LHS over all  $s$ , and maximise RHS over all  $x$  s.t.  $Ax = y$ :

$$\min_s \frac{|A^\dagger(y)s|_\infty}{\left| \sum_{j=0}^N s_j \right|} \geq \max_{x: Ax=y} \frac{1}{|x|_1} .$$



# Proof

- Hölder's inequality for vectors:  $|\langle a|b\rangle| \leq |a|_1 |b|_\infty$ .
- In our case:

$$|A^\dagger(y)s|_\infty |x|_1 \geq |x^\dagger A^\dagger(y)s| = |y^\dagger(y)s| = \left| \sum_j s_j |y_j|^2 \right| = \left| \sum_j s_j \right|$$

- Thus, for all  $s$  and for all  $x$  satisfying  $Ax = y$ :

$$\frac{|A^\dagger(y)s|_\infty}{\left| \sum_{j=0}^N s_j \right|} \geq \frac{1}{|x|_1},$$

- Minimise LHS over all  $s$ , and maximise RHS over all  $x$  s.t.  $Ax = y$ :

$$\min_s \frac{|A^\dagger(y)s|_\infty}{\left| \sum_{j=0}^N s_j \right|} \geq \max_{x: Ax=y} \frac{1}{|x|_1} = \sqrt{P_{\max}}.$$

□



# Maximal Success Probability

- We have seen that the expression for the  $a_{jk}$  depends on the transmittivity  $T$ , which is also a tunable parameter for maximising  $P$ . Hence we write  $A = A(T)$ .
- The claim is that the phase-shift gate  $c_k \mapsto y_k c_k$ , with  $y_k = 1$ , for  $k = 0 \dots N - 1$ , and  $y_N = -1$ , can be obtained with probability at most  $P_{\max} = 1/N^2$ .
- This can be proven using the Eisert bound.
- Need to find  $s(T)$  such that

$$\frac{|A(T)^\dagger(y s(T))|_\infty}{|\sum_{j=0}^N s(T)_j|} = 1/N.$$



# Maximal Success Probability

- For  $N = 2$ , and restricting to real  $T$ , the following  $s$  does the trick:

$$s_1, s_2 = \frac{1}{2} \begin{cases} 1/(1-T), & 0, & \text{if } T \in [-1, 1 - \sqrt{2}), \\ 0, & 1/(1+T^2), & \text{if } T \in [1 - \sqrt{2}, 0), \\ 1, & 1/2, & \text{if } T \in [0, 1), \end{cases}$$

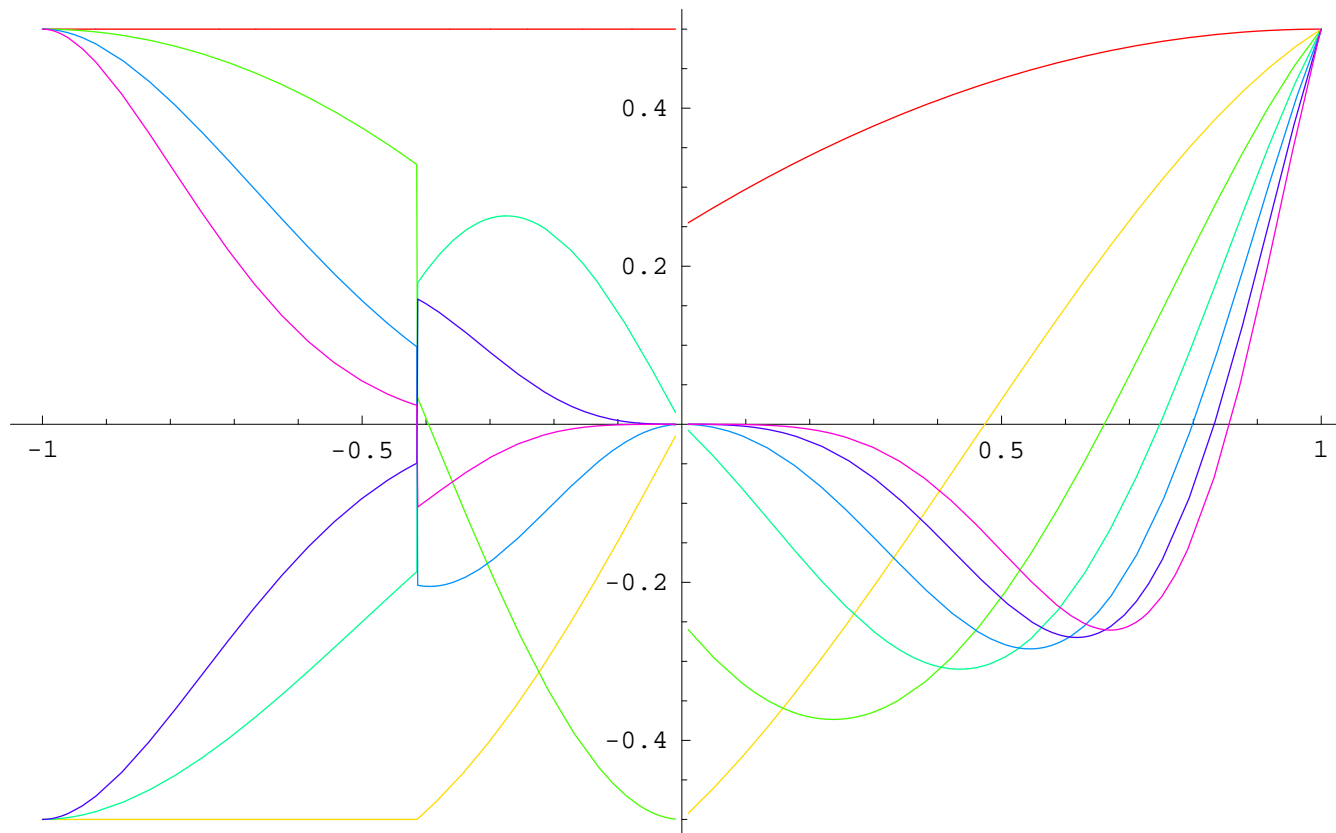
and  $s_0 = 1 - s_1 - s_2$ .

- For general  $N$ , similar choices can most likely be found (To Do!)
- On the next page I plot the first 7 entries of  $A(T)^\dagger(y_s(T))$ , showing that they are all bounded by  $1/2$ , which suggests that  $P$  is bounded by  $1/4$ .
- A real proof is of course needed that this is true for all entries (Need More Chocolate!)



# Maximal Success Probability

First 7 entries of  $A^\dagger(y_s)$  versus  $T$





## Is this bound achievable?

- To prove that the upper bound  $P \leq 1/N^2$  is achievable, consider the original problem formulation (modified with an additional maximisation over  $T$ )

$$P_{\max} = \max_T \max_x \left\{ \frac{1}{|x|_1^2} : A(T)x = y \right\}$$

- Claim **(a)**:  $P_{\max} = 1/N^2$  is achieved for  $T = 1 - 2^{1/N}$  and  $x = A_N(T)^{-1}y$  where  $A_N$  is the leftmost  $(N + 1) \times (N + 1)$  submatrix of  $A$ .
- Claim **(b)**: for this  $x$ ,  $x_l = 0$ ,  $l \geq N$ .
- We already knew that  $N + 1$  non-zero entries were enough; now this says that  $N$  non-zero entries suffice. Moreover, its the entries with lowest photon number.
- Hence, maximum  $P$  is achieved with the lowest possible ancilla dimension!



# Entries of $A$ and its Properties

- Some preliminaries...
- First we need an explicit formula for the  $a_{jk}$ .
- The permanents can be calculated explicitly, using the Lagrange expansion formula for the permanent. After a long calculation we find

$$a_{kl} = (T^*)^{l-k} \sum_{m=0}^k \binom{k}{m} \binom{l+m}{m} t^m,$$

where  $t = |T|^2 - 1$ .

- This expression does not explicitly contain  $N$ . This means that the  $(N+1) \times \infty$  matrix  $A$  is just a submatrix of the infinite-dimensional  $\mathcal{A} := (a_{kl})_{k,l=0}^{\infty}$ .



# Entries of $A$ and its Properties

- This matrix  $\mathcal{A}$  can be decomposed as a product of a number of matrices with a very simple structure.
- Define the matrices

$$D(x) := \text{diag}(1, x, x^2, \dots, x^\infty)$$
$$\Delta := \left( \binom{k}{l} \right)_{k,l=0}^\infty.$$

$\Delta$  is a representation of Pascal's triangle by a lower triangular matrix.

- A straightforward calculation leads to the decomposition

$$\mathcal{A} = D(T^*) \Delta D\left(\frac{|T|^2 - 1}{|T|^2}\right) \Delta^T D(T).$$



# Entries of $A$ and its Properties

- We will also need the matrix  $A_N$ , which is the upper left  $(N + 1) \times (N + 1)$  submatrix of  $\mathcal{A}$ . Define  $D_N(x)$  and  $\Delta_N$  accordingly.
- Since  $\det(\Delta_N) = 1$ ,  $\det(A_N) = \det(D_N(t)) = t^{N(N+1)/2}$ . Thus  $A_N$  is invertible provided  $t \neq 0$ , i.e.  $|T| < 1$ .
- The inverse of  $\Delta_N$  is  $\Delta_N$  itself, but with additional negative signs arranged in a checkerboard pattern:

$$\sum_{k=0}^N (-1)^{k-m} \binom{p}{k} \binom{k}{m} = \delta_{pm}.$$

- This easily gives the inverse of  $A_N$  as

$$(A_N^{-1})_{jk} = (-T^*)^{k-j} \sum_{l=0}^N t^{-l} \binom{l}{k} \sum_{p=0}^N \binom{p}{j} \binom{p}{l}.$$



# Equations: buy one, get one free

- We're claiming that  $x$  with  $x_l = 0$  for  $l > N$  are optimal.
- The only such  $x$  that satisfies  $Ax = y$  is  $x = A_N^{-1}y$ .
- With  $t = |T|^2 - 1$ ,

$$\begin{aligned}x_j &= \sum_{k=0}^N (A_N^{-1})_{jk} y_k = \sum_{k=0}^N (A_N^{-1})_{jk} - 2(A_N^{-1})_{jN} \\ &= (-T^*)^{-j} \left( \sum_{p=0}^{N-1} \binom{p}{j} \left(1 + \frac{1 - T^*}{t}\right)^p - \left(\frac{-T^*}{t}\right)^N \binom{N}{j} \right).\end{aligned}$$



# Equations: buy one, get one free

- In  $T = 1 - 2^{1/N}$ ,  $x_j$  simplifies to

$$x_j = (-1)^j \frac{1}{(T+1)^{N-1}} \sum_{l=0}^{N-j-1} \binom{N}{l} T^l.$$

- Thus, clearly,  $x_N = 0$  in that point: claim **(b)**.
- For  $T = 1 - 2^{1/N}$  the sum over  $l$  always yields positive values. Thus

$$|x_j| = \frac{1}{(T+1)^{N-1}} \sum_{l=0}^{N-j-1} \binom{N}{l} T^l.$$



# Equations: buy one, get one free

- Summing over all  $j$  finally yields

$$\begin{aligned}\sum_{j=0}^{N-1} |x_j| &= \frac{1}{(T+1)^{N-1}} \sum_{j=0}^{N-1} \sum_{l=0}^{N-j-1} \binom{N}{l} T^l \\ &= \frac{1}{(T+1)^{N-1}} \sum_{l=0}^{N-1} (N-l) \binom{N}{l} T^l \\ &= N.\end{aligned}$$

- Thus, for the success probability we obtain claim **(a)**:

$$P = \frac{1}{N^2}.$$

□



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# Conclusion

- We have proven that the  $N + 1$  dimensional phase-gate can be implemented using a P-A-D circuit with success probability  $1/N^2$ .
- We have shown that, for  $N = 2$ , this is the maximum.
- It remains to find trial functions  $s_1(T), \dots, s_N(T)$  to show that this is the maximum for any  $N > 2$ .
- We have shown that this success probability, taken as a function of the chosen ancilla photon numbers, can be obtained if the ancilla contains as few photons as possible (namely  $N - 1$ ).
- This is fortunate because in this case the least quantum-state engineering is needed to generate the ancilla in the preparation stage 'P' and, furthermore, decoherence affects the ancilla only minimally.