

UNIVERSITY OF LONDON

MSc/MSci EXAMINATION 2023

For Students of the University of London

DO NOT TURN OVER UNTIL TOLD TO BEGIN

**PH4211: STATISTICAL MECHANICS
PH5211: STATISTICAL MECHANICS
PH5911: STATISTICAL MECHANICS
PAPER FOR FIRST SIT/RESIT CANDIDATES**

Time Allowed: **TWO AND A HALF hours**

Answer **THREE** questions

No credit will be given for attempting any further questions

Approximate part-marks for questions are given in the right-hand margin

The total available marks add up to 120

All College-approved Calculators are permitted

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2022-23

GENERAL PHYSICAL CONSTANTS

Permeability of vacuum	μ_0	=	$4\pi \times 10^{-7}$	H m^{-1}
Permittivity of vacuum	ε_0	=	8.85×10^{-12}	F m^{-1}
	$1/4\pi\varepsilon_0$	=	9.0×10^9	m F^{-1}
Speed of light in vacuum	c	=	3.00×10^8	m s^{-1}
Elementary charge	e	=	1.60×10^{-19}	C
Electron (rest) mass	m_e	=	9.11×10^{-31}	kg
Unified atomic mass constant	m_u	=	1.66×10^{-27}	kg
Proton rest mass	m_p	=	1.67×10^{-27}	kg
Neutron rest mass	m_n	=	1.67×10^{-27}	kg
Ratio of electronic charge to mass	e/m_e	=	1.76×10^{11}	C kg⁻¹
Planck constant	h	=	6.63×10^{-34}	J s
	$\hbar = h/2\pi$	=	1.05×10^{-34}	J s
Boltzmann constant	k_B	=	1.38×10^{-23}	J K⁻¹
Stefan-Boltzmann constant	σ	=	5.67×10^{-8}	W m⁻² K⁻⁴
Gas constant	R	=	8.31	J mol⁻¹ K⁻¹
Avogadro constant	N_A	=	6.02×10^{23}	mol⁻¹
Gravitational constant	G	=	6.67×10^{-11}	N m² kg⁻²
Acceleration due to gravity	g	=	9.81	m s⁻²
Volume of one mole of an ideal gas at STP		=	2.24×10^{-2}	m³
One standard atmosphere	P_0	=	1.01×10^5	N m⁻²

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1. (a) The partition function for a single particle of mass m at temperature T , in a box of volume V is given by

$$z = V \left(\frac{mkT}{2\pi\hbar^2} \right)^{3/2} = \frac{V}{\Lambda^3}.$$

Why is the quantity Λ referred to as the thermal de Broglie wavelength? [4]

- (b) Carefully explain why the partition function for a collection of N similar but *distinguishable* objects, each with partition function z , is expressed as

$$Z = z^N.$$

[10]

- (c) On the assumption that $Z = z^N$ is correct for a gas of particles, show that the Helmholtz free energy for such a gas would be given by

$$F = -NkT \ln \left\{ V \left(\frac{mkT}{2\pi\hbar^2} \right)^{3/2} \right\}.$$

[4]

- (d) This expression gives a Helmholtz free energy that is *not extensive*. Explain what this means, and explain why this is a problem. [8]

- (e) Briefly discuss Gibbs's proposal, before the quantum era, for how this problem could be resolved. [4]

- (f) In the quantum view the partition function is

$$Z = \frac{1}{N!} z^N.$$

What is the reason for the $1/N!$ here? Explain why no problem appears with the Helmholtz free energy in this case and discuss why this resolution relies on the *thermodynamic limit*. [10]

[You may need Stirling's approximation: $\ln n! \approx n \ln(n/e)$.]

2. (a) In the Weiss model of ferromagnetism it is assumed that the magnetic moments are subject to an additional “mean” magnetic field \mathbf{b}

$$\mathbf{b} = \lambda \mathbf{M}$$

where \mathbf{M} is the magnetization and λ is a constant.

Explain briefly the origin of this field.

[8]

- (b) The magnetization of a collection of N non-interacting spin $\frac{1}{2}$ magnetic moments μ at a temperature T is given by

$$M = M_0 \tanh\left(\frac{\mu B}{k_B T}\right)$$

where M_0 is the saturation magnetization $M_0 = N\mu$ and \mathbf{M} is parallel to the applied magnetic field \mathbf{B} .

Show that the Weiss model leads to a spontaneous magnetization, in the absence of an external magnetic field, given by

$$\frac{M}{M_0} = \tanh\left(\frac{M T_c}{M_0 T}\right)$$

where $T_c = \lambda M_0^2 / N k_B$.

[8]

- (c) i. Sketch the behaviour of the spontaneous magnetization as a function of temperature. [4]
 ii. What is the interpretation of T_c ? [2]
 iii. Discuss the order of the transition. [2]
- (d) When T is very close to T_c then M/M_0 is very small. By expanding the hyperbolic tangent ($\tanh x \approx x - x^3/3 + \dots$) show that, in the vicinity of T_c , the magnetization behaves as

$$\frac{M}{M_0} \propto \left(1 - \frac{T}{T_c}\right)^{1/2}.$$

[8]

- (e) The critical exponent β describes the behaviour of the order parameter in the vicinity of the critical point. How is β defined?
 How does the result in (d) compare with the behaviour of real systems?
 Discuss the difference.

[8]

3. In the Landau theory of phase transitions the free energy is expanded in powers of the order parameter.

(a) In the context of phase transitions, what is meant by the term *order parameter*? [4]

(b) Explain why only a limited number of terms are needed. What determines the highest power in the expansion? [6]

(c) What is the distinction between a *conserved* order parameter and a *non-conserved* order parameter? Give an example of each. Describe how the equilibrium state of the system is determined in these two cases. [7]

(d) What is the order parameter of a binary alloy? Sketch the free energy for a temperature *above* and *below* the critical temperature. [6]

(e) The free energy of mixing for the binary alloy may be written

$$F_m = 2NkT_c x(1-x) + NkT [x \ln x + (1-x) \ln(1-x)]$$

where x is the order parameter.

Explain the origin of the two terms in this expression. [6]

(f) The expansion of F_m may be written as

$$F_m = F_0 + 2Nk \left\{ (T - T_c) \left(x - \frac{1}{2}\right)^2 + \frac{2}{3}T_c \left(x - \frac{1}{2}\right)^4 + \frac{16}{15}T_c \left(x - \frac{1}{2}\right)^6 + \dots \right.$$

Discuss the Landau truncation of this expression; in particular explain at which term the series should be terminated. Why is the expansion written in powers of $x - \frac{1}{2}$? [7]

(g) Within this model the order parameter critical exponent β has the value $\frac{1}{2}$. Show how this follows from the Landau free energy. [4]

4. Write an essay on the logical foundations of Statistical Mechanics and its connection with Thermodynamics. You should include a discussion of the nature of *equilibrium states*, the emergence of the concept of *temperature*, and the importance of the *thermodynamic limit*. [40]

5. The force on a Brownian particle in one dimension may be written as $F(t) = f(t) - v/\mu$ where $f(t)$ is the randomly fluctuating force experienced by the particle, v is the velocity and μ the mobility of the particle.

(a) Discuss the separation of the force into the two parts. In particular, explain qualitatively how the damping force, proportional to the velocity, arises. [10]

(b) Show that the equation of motion for the Brownian particle of mass M (the Langevin equation) may be written as [4]

$$M \frac{dv(t)}{dt} + \frac{1}{\mu} v(t) = f(t).$$

(c) The solution to the Langevin equation is given by

$$v(t) = v(0)e^{-t/M\mu} + \frac{1}{M} \int_0^t e^{(s-t)/M\mu} f(s) ds.$$

i. Show, using appropriate approximations, that the equilibrium mean square velocity may be expressed as $\langle v^2 \rangle = \frac{\mu}{2M} \int_{-\infty}^{\infty} \langle f(0)f(t) \rangle dt$. [8]

ii. By invoking equipartition, show it follows that the mobility may be expressed $\frac{1}{\mu} = \frac{1}{2kT} \int_{-\infty}^{\infty} \langle f(0)f(t) \rangle dt$. [3]

(d) Discuss how this may be regarded as an example of the *fluctuation-dissipation theorem*. [5]

(e) The voltage across an $L - R$ circuit is given in terms of the current $I(t)$ by

$$L \frac{dI(t)}{dt} + RI(t) = V(t)$$

where $V(t)$ may be regarded as a fluctuating voltage. By making the appropriate identifications, show that the resistance may be related to the voltage fluctuations through [6]

$$R = \frac{1}{2kT} \int_{-\infty}^{\infty} \langle V(0)V(t) \rangle dt.$$

(f) How may such voltage fluctuations be observed? [4]

END